

3rd International Conference Advanced Mechanics: Structure, Materials, Tribology

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AIPCP25-CF-AMSMT2025-00038 | Article

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Application of the Mixed Approximation for Solving Two Problems of the Theory of Elasticity by the Finite Element Method

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Annotation: The article formulates a discrete problem of the theory of elasticity using the mixed finite element method. Finite elements, a set of nodes, approximation domains, and a series of finite-dimensional subspaces to which the solutions of the discrete problem must belong are defined. An algorithm is provided for constructing piecewise polynomial functions for both displacements and stresses (strains). The existence and uniqueness of the solution to the discrete problem are proven.

Keywords: finite functions, triangulation, subspaces, mixed finite element method late

INTRODUCTION.

The mixed finite element method (MFEM) is based on the idea of simultaneously approximating displacements, stresses, and strains. This leads to an increase in the number of corresponding equations and a higher computational cost. Therefore, it becomes necessary to conduct studies aimed at exploring the capabilities of MFEM that is, the investigation of the existence, uniqueness, and stability of solutions, as well as assessing the method's effectiveness in solving specific problems in mechanics. The present work is devoted to the development of MFEM in the context of linear elasticity problems and its mathematical justification. Conditions for the solvability of the corresponding discrete problems are obtained [1-6].

STATEMENT OF THE PROBLEMS (ISSUE)

Let us consider a boundary value problem of linear elasticity in a generalized formulation [7].

$$(\varepsilon, \chi)_X = (Bu, \chi)_X, \quad \forall \chi \in X \quad (1)$$

$$(\sigma, \tau)_X = (D\varepsilon, \tau)_X - (D\zeta, \tau)_X, \quad \forall \tau \in X \quad (2)$$

$$(\sigma, Bv)_X = \langle \rho, v \rangle_u, \quad \forall v \in U \quad (3)$$

To construct mixed finite element methods (MFEM), we perform an approximation, replacing the spaces U and X , where $u \in U$ and $\sigma, \varepsilon \in X$, with finite-dimensional subspaces $U_h \subset U$ and $X_h \subset X$. In this case, the parameter is $h \rightarrow 0$.

On the domain Ω , we perform the discretization Γ_h :

$$\Gamma_h = \left\{ \bigcup_{m=1, \dots, N_m(h)} \bar{\Omega}_m = \bar{\Omega}; \Omega_m \cap \Omega_l = \emptyset, \forall m \neq l; \right. \\ \left. \max_{m=1, \dots, N_m(h)} h_m \leq h; h_m \leq c s_m, \forall \Omega_m \right\} \quad (4)$$

Here, the diameter of the circumscribed sphere around Ω_m is h_m , and the diameter of the inscribed sphere in Ω_m is denoted by s_m . The number of subsets $N_m(h)$ is denoted by c — a constant that does not depend on h .

Let the sets of nodes for displacements and stresses (strains) be denoted by ε_h and Y_h , respectively [8-9]:

$$\varepsilon_h = (X_\alpha)_{\alpha=1}^{N_\alpha(h)}, Y_h = (X_\gamma)_{\gamma=1}^{N_\gamma(h)}. \quad (5)$$

The sets of piecewise polynomial functions for $x_\alpha \in \varepsilon_h$ and $x_\gamma \in Y_h$ are defined as follows: $(\varphi_{h\alpha}(x))_{\alpha=1}^{N_\alpha(h)}$ and $(\psi_{h\gamma}(x))_{\gamma=1}^{N_\gamma(h)}$, possessing certain properties.

$$\varphi_{h\alpha}(x_\beta) = \delta_{\alpha\beta}, \quad \forall x_\alpha, x_\beta \in \varepsilon_h \quad (6)$$

$$\psi_{h\gamma}(x_\mu) = \delta_{\gamma\mu}, \quad \forall x_\gamma, x_\mu \in Y_h \quad (7)$$

$$\text{supp } \varphi_{h\alpha}(x) = \{x | x \in \Lambda_{h\alpha}\} \quad (8)$$

$$\text{supp } \psi_{h\gamma}(x) = \{x | x \in \Lambda_{h\gamma}\} \quad (9)$$

the collection of subsets in Ω_m containing the nodes $x_\alpha \in \varepsilon_h$ and $x_\gamma \in Y_h$ is denoted by $\Lambda_{h\alpha}$ and $\Lambda_{h\gamma}$, respectively.

Using linear combinations of the above-mentioned basis functions, we define the functions for displacements and stresses (strains):

$$v_{hi}(x) = \sum_{x_\alpha \in \varepsilon_h} v_{hi}(x_\alpha) \varphi_{h\alpha}(x), \quad \forall x \in \Omega \quad (10)$$

$$\tau_{hij}(x) = \sum_{x_\gamma \in Y_h} \tau_{hij}(x_\gamma) \psi_{h\gamma}(x), \quad \forall x \in \Omega \quad (11)$$

They, in turn, are $\Phi_h(\Omega)$ and $\Psi_h(\Omega)$, which are finite element spaces satisfying the following conditions:

$$\Phi_h(\Omega) \subset H^1(\Omega), \psi_h(\Omega) \subset L_2(\Omega), \quad (12)$$

$$U_h = \{v_h | v_h = (v_{hi})_{i=1}^n, v_{hi} \in \Phi_h(\Omega), v_{hi}|_{\Gamma_u} = 0\} \quad (13)$$

$$X_h = \{\tau_h | \tau_h = (\tau_{hij})_{i,j=1}^n, \tau_{hij} \in \Psi_h(\Omega)\} \quad (14)$$

Therefore, (13) and (14) are the approximating spaces for displacements and stresses (strains), respectively, with the following conditions holding: $U_h \subset U$ and $X_h \subset X$.

Let us consider the finite-dimensional spaces V_h and Z_h . Their elements are the following basis functions:

$$\hat{\mathcal{G}}_h = (\hat{\mathcal{G}}_{h\alpha})_{\alpha=1}^{N_\alpha(h)} \text{ and } \hat{\tau}_h = (\hat{\tau}_{h\gamma})_{\gamma=1}^{N_\gamma(h)}, \text{ with } \hat{\mathcal{G}}_{h\alpha} = (\hat{\mathcal{G}}_{hia})_{i=1}^n \text{ and } \hat{\tau}_{h\gamma} = (\tau_{hij\gamma})_{i,j=1}^n.$$

It is assumed that $\hat{\mathcal{G}}_h \in V_h$ and $\hat{\tau}_h \in Z_h$ hold if $P_h \hat{\mathcal{G}}_h \in U_h$ and $R_h \hat{\tau}_h \in X_h$. Then, P_h and R_h are mapping operators from V_h to U_h and from Z_h to X_h , respectively. They are defined as follows:

$$\mathcal{G}_{hi}(x) = (P_h \hat{\mathcal{G}}_h)_i(x) = \sum_{x_\alpha \in \varepsilon_h} \mathcal{G}_{hia} \varphi_{h\alpha}(x), \quad \forall \hat{\mathcal{G}}_h \in V_h \quad (15)$$

$$\tau_{hij}(x) = (R_h \hat{\tau}_h)_{ij}(x) = \sum_{x_\gamma \in Y_h} \tau_{hij\gamma} \psi_{h\gamma}(x), \quad \forall \hat{\tau}_h \in Z_h \quad (16)$$

From conditions (6) and (7), we obtain:

$$\mathcal{G}_{hi}(x_\alpha) = \mathcal{G}_{hi\alpha}, \quad \forall x_\alpha \in \mathcal{E}_h \quad (17)$$

$$\tau_{hij}(x_\gamma) = \tau_{hij\gamma}, \quad \forall x_\gamma \in Y_h \quad (18)$$

Thus, the following functions correspond to the numerical parameters (17) and (18):

$$\mathcal{G}_h = P_h \hat{\mathcal{G}}_h, \quad \forall \hat{\mathcal{G}}_h \in V_h \quad (19)$$

$$\tau_h = R_h \hat{\tau}_h, \quad \forall \hat{\tau}_h \in Z_h \quad (20)$$

With the following formulas, we introduce a set of functions: $(\mu_{hi\alpha}(x))_{\alpha=1}^{N_\alpha(h)} (i=1, \dots, n)$.

where

$$\mu_{hi\alpha}(x) = \frac{1}{2} \frac{\partial \phi_{hi\alpha}(x)}{\partial x_i} = \begin{cases} \frac{1}{2} \frac{\partial \lambda_{hi\alpha}^{(m)}(x)}{\partial x_i}, & \forall x \in \Lambda_{hi\alpha} \\ 0, & \forall x \notin \Lambda_{hi\alpha} \end{cases} \quad (21)$$

The notation $\lambda_{hi\alpha}^{(m)}(x)$ refers to polynomial coordinate functions within the bounds of Ω_m , satisfying the condition:

$\lambda_{hi\alpha}^{(m)}(x_\beta) = \delta_{\alpha\beta}$. Let Y_h denote the linear span of the functions $(\mu_{hi\alpha}(x))_{\alpha=1}^{N_\alpha(h)}$:

$$Y_h = \{ \bar{\tau}_h | \bar{\tau}_h = (\bar{\tau}_{hij})_{i,j=1}^n, \bar{\tau}_{hij} = \bar{\tau}_{hij}, \bar{\tau}_{hij}(x) = \sum_{x_\alpha \in \mathcal{E}_h} v_{hi\alpha} \mu_{hj\alpha}(x) + v_{hj\alpha} \mu_{hi\alpha}(x), \forall \hat{v} \in V_h \} \quad (22)$$

Based on the independence of $(\varphi_{hi\alpha}(x))_{\alpha=1}^{N_\alpha(h)}$, we obtain the independence of the functions $(\mu_{hi\alpha}(x))_{\alpha=1}^{N_\alpha(h)}$.

Therefore, the dependence of $Y_h = BU_h$ is interpreted by Y_h as a Hilbert space dependence. Then the linear operator $B_h = BP_h$ acts from U_h to Y_h , defined by the following relations:

$$(B_h \hat{v}_h)_{ij} = \sum_{x_\alpha \in \mathcal{E}_h} v_{hi\alpha} \mu_{hj\alpha}(x) + v_{hj\alpha} \mu_{hi\alpha}(x) \quad (23)$$

Y_h is the image of the operator B_h .

SOLUTION OF THE PROBLEM

Find $(u_h, \sigma_h, \varepsilon_h) \in U_h \times X_h \times X_h$ such that

$$(\varepsilon_h, \chi_h)_X = (Bu_h, \chi_h)_X, \quad \forall \chi_h \in X_h; \quad (24)$$

$$(\sigma_h, \tau_h)_X = (\wp \varepsilon_h, \tau_h)_X, \quad \forall \tau_h \in X_h; \quad (25)$$

$$(\sigma_h, Bv_h)_X = \langle \rho, v_h \rangle_U, \quad \forall v_h \in U_h; \quad (26)$$

An arbitrary function from U_h and X_h is uniquely represented in the form of (19) and (20). Then (24)–(26) can be represented in the following form:

$$(R_h, R_h, \hat{\chi}_h)_X = (B_h \hat{u}_h, R_h \hat{\chi}_h)_X, \quad (27)$$

$$(R_h, R_h, \hat{\tau}_h)_X = (\wp R_h \hat{\varepsilon}_h - \wp \zeta, R_h \hat{\tau}_h)_X, \quad (28)$$

$$(R_h, R_h, \hat{v}_h)_X = \langle \rho, P_h \hat{v}_h \rangle_U, \quad (29)$$

The adjoint operators P_h^* , R_h^* and B_h^* to the operators P_h , R_h and B_h are defined by the relations:

$$\hat{\mathcal{G}}_h \rightarrow \langle g, P_h \hat{\mathcal{G}}_h \rangle_U = \langle P_h^* g, \hat{\mathcal{G}}_h \rangle_{V_h}; \quad \forall g \in U^*, \forall \hat{\mathcal{G}}_h \in V_h; \quad (30)$$

$$\hat{\chi}_h \rightarrow (\tau, R_h \hat{\chi}_h)_X = \langle R_h^* \tau, \hat{\chi}_h \rangle_{Z_h}; \quad \forall \tau \in X^*, \forall \hat{\chi}_h \in Z_h; \quad (31)$$

$$\hat{\mathcal{G}}_h \rightarrow (\tau, B_h \hat{\mathcal{G}}_h)_X = \langle B_h^* \tau, \hat{\mathcal{G}}_h \rangle_{V_h}; \quad \forall \tau \in X^*, \forall \hat{\mathcal{G}}_h \in V_h. \quad (32)$$

Next, we introduce the linear operator M_h :

$$\hat{\chi}_h \rightarrow (R_h \hat{\tau}, R_h \hat{\chi}_h)_X = \langle M_h \hat{\tau}, \hat{\chi}_h \rangle_{Z_h}; \quad \forall \hat{\tau}_h, \hat{\chi}_h \in Z_h; \quad (33)$$

Then $M_h : Z_h \rightarrow Z_h^*$ is represented in the form:

$$M_h = R_h^* R_h \quad \text{B} \quad Z_h^* \quad (34)$$

In detail, we define the linear operator M_h as G_h :

$$\hat{\chi}_h \rightarrow (\wp R_h \hat{\tau}, R_h \hat{\chi}_h)_X = \langle G_h \hat{\tau}, \hat{\chi}_h \rangle_{Z_h}; \quad \forall \hat{\tau}_h, \hat{\chi}_h \in Z_h; \quad (35)$$

where $G_h : Z_h \rightarrow Z_h^*$ is represented in the form:

$$G_h = R_h^* \wp R_h \quad \text{B} \quad Z_h^* \quad (36)$$

Next, we define the linear operator H_h :

$$\hat{\chi}_h \rightarrow (B_h \hat{v}_h, R_h \hat{\chi}_h)_X = \langle H_h \hat{v}_h, \hat{\chi}_h \rangle_{Z_h}; \quad \forall \hat{v}_h \in V_h, \forall \hat{\chi}_h \in Z_h; \quad (37)$$

where $H_h : V_h \rightarrow Z_h^*$ is represented in the form:

$$H_h = R_h^* B_h \quad \text{B} \quad Z_h^* \quad (38)$$

Let H_h' denote the transpose of the operator H_h , where $H_h' : Z_h \rightarrow V_h^*$ is defined as follows:

$$\hat{v}_h \rightarrow (\hat{\tau}_h, H_h' \hat{v}_h)_{Z_h} = \langle H_h' \hat{\tau}_h, \hat{v}_h \rangle_{V_h}; \quad \forall \hat{\tau}_h \in Z_h, \forall \hat{v}_h \in V_h; \quad (39)$$

Then H_h' can be represented in the form:

$$H_h' = B_h' R_h \quad \text{B} \quad Z_h^* \quad (40)$$

Based on the definition of the operators M_h , H_h , G_h , H_h' , equations (29)–(31) can be written in the form of operator equations:

$$M_h \hat{\varepsilon}_h = H_h \hat{u}_h \quad \text{B} \quad Z_h^* \quad (41)$$

$$M_h \hat{\sigma}_h = G_h \hat{\varepsilon}_h - \hat{\xi}_h \quad \text{B} \quad Z_h^* \quad (42)$$

$$H_h' \hat{\sigma}_h = \hat{\rho}_h \quad \text{B} \quad V_h^* \quad (43)$$

In this case,

$$\hat{\xi}_h = R_h^* D \zeta \in Z_h^* \quad (44)$$

$$\hat{\rho}_h = P_h^* \rho \in V_h^* \quad (45)$$

Taking into account that G_h can be represented in the form:

$$G_h = M_h \hat{D}_h \quad \text{B} \quad Z_h^* \quad (46)$$

where \hat{D}_h is the restriction of the operator D to Z_h , and the element $\hat{\xi}_h \in Z_h^*$ is defined by the formula:

$$\hat{\xi}_h = M_h \hat{D}_h \hat{\zeta}_h \in Z_h^* \quad (47)$$

The element $\hat{\zeta}_h \in Z_h$ is the restriction of $\zeta \in X$ to Y_h .

Taking into account the above formulas, the system of equations (41)–(43) can be written in the form:

$$M_h \hat{\varepsilon}_h = H_h \hat{u}_h \quad \text{B} \quad Z_h^* \quad (48)$$

$$\hat{\sigma}_h = \hat{D}_h (\hat{\varepsilon}_h - \hat{\xi}_h) \quad \text{B} \quad Z_h^* \quad (49)$$

$$H_h' \hat{\sigma}_h = \hat{\rho}_h \quad \text{B} \quad V_h^* \quad (50)$$

The latter system can be represented as a single operator equation:

$$A_h \hat{u}_h = \hat{f}_h \quad \text{B} \quad V_h^* \quad (51)$$

In this case, the linear operator $A_h : V_h \rightarrow V_h^*$ is defined by the formula:

$$A_h = H_h' \hat{D}_h M_h^{-1} H_h \quad \text{B} \quad V_h^* \quad (52)$$

The element $\hat{f}_h \in V_h^*$ is defined by the formula: ой:

$$\hat{f}_h = \hat{\rho}_h + H_h' \hat{D}_h \hat{\zeta}_h \in V_h^* \quad \text{B} \quad V_h^* \quad (53)$$

Thus, equations (48)–(50) define the system of finite element equations for the problem of elasticity theory [2].

Let's consider the question of existence and uniqueness of the approximate solution. To this end, we define scalar products on the spaces Z_h and V_h .

$$\widehat{f}_h = \widehat{\rho}_h + H_h^* \widehat{D}_h \widehat{\zeta}_h \in V_h^* (\widehat{\tau}_h, \widehat{\chi}_h)_{Z_h} = (R_h \widehat{\tau}_h, R_h \widehat{\chi}_h)_X; \quad \forall \widehat{\tau}_h, \widehat{\chi}_h \in Z_h \quad \text{B} \quad V_h^* \quad (54)$$

$$\|\cdot\|_{Z_h} : \widehat{\tau}_h \in Z_h \rightarrow \|\widehat{\tau}_h\|_{Z_h} = (\widehat{\tau}_h, \widehat{\tau}_h)_{Z_h}^{1/2} \quad \text{B} \quad V_h^* \quad (55)$$

It follows that (55) is a norm on Z_h and that it is a Hilbert space isomorphic to X_h . According to (54) and the operator $M_h : Z_h \rightarrow Z_h^*$, we obtain:

$$(\widehat{\tau}_h, \widehat{\chi}_h)_{Z_h} = \langle M_h \widehat{\tau}_h, \widehat{\chi}_h \rangle_{Z_h} = \langle \widehat{\tau}_h, M_h \widehat{\chi}_h \rangle_{Z_h}; \quad \forall \widehat{\tau}_h, \widehat{\chi}_h \in Z_h \quad \text{B} \quad V_h^* \quad (56)$$

And this ensures that M_h is the operator of canonical isometry from Z_h into Z_h^* ; then there exists M_h^{-1} , the inverse of M_h , called the Ross operator, and $M_h^{-1} : Z_h^* \rightarrow Z_h$ is unique. Defining in Z_h^* the scalar product:

$$(\widehat{\xi}_h, \widehat{\eta}_h)_{Z_h^*} = (M_h^{-1} \widehat{\xi}_h, M_h^{-1} \widehat{\eta}_h)_{Z_h} = \langle M_h^{-1} \widehat{\xi}_h, \widehat{\eta}_h \rangle_{Z_h} = \langle \widehat{\xi}_h, M_h^{-1} \widehat{\eta}_h \rangle_{Z_h}; \quad \forall \widehat{\xi}_h, \widehat{\eta}_h \in Z_h^* \quad (57)$$

Let us establish that Z_h^* is also a Hilbert space such that.

$$\|\cdot\|_{Z_h^*} : \widehat{\xi}_h \in Z_h^* \rightarrow \|\widehat{\xi}_h\|_{Z_h^*} = (\widehat{\xi}_h, \widehat{\xi}_h)_{Z_h^*}^{1/2} \quad (58)$$

Now we construct a scalar product on V_h ; for this we establish a correspondence between X_h and Y_h . To each element $B\mathcal{G}_h \in Y_h$, $\forall \mathcal{G}_h \in U_h$ we associate its projection onto X_h . This association is an operator in X . Denote the operator that performs this by J_h . By definition J_h is the operator of orthogonal projection of Y_h onto X_h . The element $J_h B\mathcal{G}_h$ is the orthogonal projection of $B\mathcal{G}_h \in Y_h$, $\forall \mathcal{G}_h \in U_h$ onto X_h . This means that the difference $J_h B\mathcal{G}_h - B\mathcal{G}_h$ is orthogonal (with respect to the inner product in) X to any arbitrary element from X_h , i.e.

$$(J_h B\mathcal{V}_h - B\mathcal{V}_h, \chi_h)_X = 0, \quad \forall \chi_h \in X_h \quad (59)$$

Based on the definitions of the operators J_h and R_h , we have:

$$(R_h \widehat{\tau}_h, R_h \widehat{\chi}_h)_X = (B_h \widehat{\tau}_h, R_h \widehat{\chi}_h)_X, \quad \forall \widehat{\chi}_h \in Z_h \quad (61)$$

As is known, $M_h : Z_h \rightarrow Z_h^*$ and $H_h : V_h \rightarrow Z_h^*$, we have:

$$M_h \widehat{\tau}_h = H_h \widehat{\tau}_h \quad \text{B} \quad Z_h^* \quad (62)$$

Thus, $\widehat{\tau}_h \in Z_h$ is uniquely determined by the formula:

$$\widehat{\tau}_h = M_h^{-1} H_h \widehat{\tau}_h \quad (63)$$

This makes it possible to define the operator J_h using the following expression:

$$\widehat{\tau}_h = J_h B\mathcal{V}_h = R_h M_h^{-1} H_h \widehat{\tau}_h \quad (64)$$

Let us note some properties of the orthoprojector J_h . Based on (59), we have:

$$(J_h B\mathcal{V}_h, B\mathcal{W}_h)_X = (J_h B\mathcal{V}_h, J_h B\mathcal{W}_h)_X = (B\mathcal{V}_h, J_h B\mathcal{W}_h)_X, \quad \forall \mathcal{V}_h, \mathcal{W}_h \in U_h \quad (65)$$

Next,

$$J_h^2 B\mathcal{V}_h = J_h (J_h B\mathcal{V}_h) = J_h B\mathcal{V}_h \quad \text{B} \quad X_h \quad (66)$$

Using (65) as well as the Cauchy–Bunyakovsky–Schwarz inequality, we find:

$$\|J_h B\mathcal{V}_h\|_X^2 = (J_h B\mathcal{V}_h, B\mathcal{V}_h)_X \leq \|J_h B\mathcal{V}_h\|_X \|B\mathcal{V}_h\|_X, \quad \forall \mathcal{V}_h \in U_h \quad (67)$$

it follows that

$$\|J_h B\mathcal{V}_h\|_X \leq \|B\mathcal{V}_h\|_X, \quad \forall \mathcal{V}_h \in U_h \quad (68)$$

the equality holds only on the common elements of X_h and Y_h . Based on (68), the norm of the operator J_h is:

$$\|J_h\|_{L(Y_h; X_h)} = \sup_{\mathcal{V}_h \in U_h / \{0\}} \frac{\|J_h B\mathcal{V}_h\|_X}{\|B\mathcal{V}_h\|_X} = 1 \quad (69)$$

Suppose

$$d_0 \leq d(h) = \inf_{\mathcal{V}_h \in U_h / \{0\}} \frac{\|J_h B\mathcal{V}_h\|}{\|B\mathcal{V}_h\|} \quad (70)$$

Here, d_0 is a positive number independent of h . The given relations make it possible to define the scalar product on V_h :

$$\begin{aligned} (\hat{v}_h, \hat{w}_h)_{V_h} &= (J_h B_h \hat{v}_h, J_h B_h \hat{w}_h)_X = (J_h B_h \hat{v}_h, B_h \hat{w}_h)_X = (B_h \hat{v}_h, J_h B_h \hat{w}_h)_X = \\ &= \langle M_h^{-1} H_h \hat{v}_h, H_h \hat{w}_h \rangle_{Z_h} = \langle H_h \hat{v}_h, M_h^{-1} H_h \hat{w}_h \rangle_{Z_h}; \quad \forall \hat{v}_h, \hat{w}_h \in V_h \end{aligned} \quad (71)$$

the norm in V_h is defined by the formula:

$$\|\cdot\|_{V_h} : \hat{v}_h \in V_h \rightarrow \|\hat{v}_h\|_{V_h} = (\hat{v}_h, \hat{v}_h)_{V_h}^{1/2} \quad (72)$$

The norm in V_h^* is defined by the formula:

$$\|\cdot\|_{V_h^*} : \hat{f}_h \in V_h^* \rightarrow \|\hat{f}_h\|_{V_h^*} = \sup_{\hat{v}_h \in V_h \setminus \{0\}} \frac{|\langle \hat{f}_h, \hat{v}_h \rangle_{V_h^*}|}{\|\hat{v}_h\|_{V_h}} \quad (73)$$

Based on the following relations:

$$\begin{aligned} (\hat{v}_h, \hat{w}_h)_{V_h} &= \langle J_h \hat{v}_h, \hat{w}_h \rangle_{V_h} = \langle H_h' M_h^{-1} H_h \hat{v}_h, \hat{w}_h \rangle_{V_h} = \langle \hat{v}_h, H_h' M_h^{-1} H_h \hat{w}_h \rangle_{V_h} = \\ &= \langle \hat{v}_h, J_h \hat{w}_h \rangle_{V_h}; \quad \forall \hat{v}_h, \hat{w}_h \in V_h \end{aligned} \quad (74)$$

It can be concluded that J_h is the operator of the canonical isometry from V_h to V_h^* and is represented as:

$$J_h = H_h' M_h^{-1} H_h \quad (75)$$

Obviously, $J_h^{-1} : V_h^* \rightarrow V_h$ is the Riesz operator for V_h , and a scalar product can be introduced on V_h^* :

$$(\hat{f}_h, \hat{g}_h)_{V_h^*} = (J_h^{-1} \hat{f}_h, J_h^{-1} \hat{g}_h)_{V_h} = \langle J_h^{-1} \hat{f}_h, \hat{g}_h \rangle_{V_h} = \langle \hat{f}_h, J_h \hat{g}_h \rangle_{V_h}; \quad \forall \hat{f}_h, \hat{g}_h \in V_h^* \quad (76)$$

and V_h^* is also a Hilbert operator/space. Based on the above relations, we will show that (51) has a unique solution. It is necessary to assume certain properties regarding the operator \hat{D}_h . We consider \hat{D}_h to be a linear, self-adjoint, positive definite, and bounded operator from Z_h to Z_h^* , with the condition that:

$$\exists \delta_0 = \text{const} > 0 : \forall \hat{\tau}_h \in Z_h, (\hat{D}_h \hat{\tau}_h, \hat{\tau}_h)_{Z_h} \geq \delta_0 \|\hat{\tau}_h\|_{Z_h}^2 \quad (77)$$

$$\exists \Delta_0 = \text{const} < \infty : \forall \hat{\tau}_h \in Z_h, \|\hat{D}_h \hat{\tau}_h\|_{Z_h} \leq \Delta_0 \|\hat{\tau}_h\|_{Z_h} \quad (78)$$

Taking into account that,

$$\langle A_h \hat{v}_h, \hat{w}_h \rangle_{V_h} = (D_h M_h^{-1} H_h \hat{v}_h, M_h^{-1} H_h \hat{w}_h)_{V_h}; \quad \forall \hat{v}_h, \hat{w}_h \in V_h \quad (79)$$

and

$$\|\hat{v}_h\|_{V_h}^2 = \langle J_h \hat{v}_h, \hat{v}_h \rangle_{V_h} = \|M_h^{-1} H_h \hat{v}_h\|_{Z_h}^2; \quad \forall \hat{v}_h \in V_h \quad (80)$$

based on (79), we have:

$$\langle A_h \hat{v}_h, \hat{v}_h \rangle_{V_h} \geq \delta_0 \|M_h^{-1} H_h \hat{v}_h\|_{Z_h}^2 = \delta_0 \|\hat{v}_h\|_{V_h}^2; \quad \forall \hat{v}_h \in V_h \quad (81)$$

Moreover, based on the Cauchy–Bunyakovsky–Schwarz inequality and (78), we obtain:

$$\|A_h \hat{v}_h\|_{V_h^*} \leq \Delta_0 \|M_h^{-1} H_h \hat{v}_h\|_{Z_h} = \Delta_0 \|\hat{v}_h\|_{V_h}; \quad \forall \hat{v}_h \in V_h \quad (82)$$

Thus, A_h is a positive definite and bounded operator V_h from V_h^* to.

Then there exists an operator A_h^{-1} acting from V_h^* to V_h such that

$$\|A_h^{-1}\|_{L(V_h^*, V_h)} \leq \frac{1}{\delta_0} \quad (83)$$

This means that equation (51) has a unique solution that depends continuously on the right-hand side, i.e., on the element $\hat{f}_h \in V_h^*$, and we obtain the following estimates:

$$\|\hat{u}_h\|_{V_h} = \|A_h^{-1} \hat{f}_h\|_{V_h} \leq \|A_h^{-1}\|_{L(V_h^*, V_h)} \|\hat{f}_h\|_{V_h^*} \leq \frac{1}{\delta_0} \|\hat{f}_h\|_{V_h^*} \quad (84)$$

$$\|\hat{\varepsilon}_h\|_{Z_h} = \|M_h^{-1} H_h \hat{u}_h\|_{Z_h} \leq \|\hat{u}_h\|_{V_h} \leq \frac{1}{\delta_0} \|\hat{f}_h\|_{V_h^*} \quad (85)$$

$$\|\hat{\sigma}_h\|_{Z_h} = \|\hat{D}_h (\hat{\varepsilon}_h - \hat{\zeta}_h)\|_{Z_h} \leq \Delta_0 \left(\|\hat{\varepsilon}_h\|_{Z_h} + \|\hat{\zeta}_h\|_{Z_h} \right) \leq \frac{\Delta_0}{\delta_0} \|\hat{f}_h\|_{V_h^*} + \Delta_0 \|\hat{\zeta}_h\|_{Z_h} \quad (86)$$

As a result, we obtain the following estimates:

$$\|u_h\|_U \leq \frac{1}{\delta_0 d_0} \|\hat{f}_h\|_{V_h^*} \quad (87)$$

$$\|\varepsilon_h\|_X \leq \frac{1}{\delta_0} \|\hat{f}_h\|_{V_h^*} \quad (88)$$

$$\|\sigma_h\|_X \leq \frac{\Delta_0}{\delta_0} \|\hat{f}_h\|_{V_h^*} + \Delta_0 \|\hat{\zeta}_h\|_{Z_h} \quad (89)$$

Since $U_h \subset U$ has this property, then by the embedding theorem and the equivalence of norms

$$\|u_h\|_{[L_2(\Omega)]^n} \leq c \|u_h\|_{[H^1(\Omega)]^n} \leq c_0 \|u_h\|_U \leq \frac{c_0}{\delta_0 d_0} \|\hat{f}_h\|_{V_h^*} \quad (90)$$

Based on the formula for determining the element $\hat{f}_h \in V_h^*$ and the norm of $\|\cdot\|_{V_h^*}$, we obtain:

$$\|\hat{f}_h\|_{V_h^*} \leq \|\hat{\rho}_h\|_{V_h^*} + \Delta_0 \|\hat{\zeta}_h\|_{Z_h} \quad (91)$$

As a result, estimates (84)–(86) take the following form:

$$\|\hat{u}_h\|_{V_h} \leq \frac{1}{\delta_0} \|\hat{\rho}_h\|_{V_h^*} + \frac{\Delta_0}{\delta_0} \|\hat{\zeta}_h\|_{Z_h} \quad (92)$$

$$\|\hat{\varepsilon}_h\|_{Z_h} \leq \frac{1}{\delta_0} \|\hat{\rho}_h\|_{V_h^*} + \frac{\Delta_0}{\delta_0} \|\hat{\zeta}_h\|_{Z_h} \quad (93)$$

$$\|\hat{\sigma}_h\|_{Z_h} \leq \frac{\Delta_0}{\delta_0} \|\hat{\rho}_h\|_{V_h^*} + \Delta_0 \left(1 + \frac{\Delta_0}{\delta_0} \right) \|\hat{\zeta}_h\|_{Z_h} \quad (94)$$

Therefore, the solution of the system of equations (48)–(50) exists and depends continuously on $(\hat{\rho}_h, \hat{\zeta}_h) \in V_h^* \times Z_h$.

It is not difficult to show that the problem (48)–(50) has a unique solution. Let $(\hat{u}_h^{(1)}, \hat{\sigma}_h^{(1)}, \hat{\varepsilon}_h^{(1)}) \in V_h \times Z_h \times Z_h$ and $(\hat{u}_h^{(2)}, \hat{\sigma}_h^{(2)}, \hat{\varepsilon}_h^{(2)}) \in V_h \times Z_h \times Z_h$ be two solutions. Substituting them into (48)–(50) and subtracting one set of equations from the other, we obtain:

$$\begin{cases} M_h (\hat{\varepsilon}_h^{(1)} - \hat{\varepsilon}_h^{(2)}) = H_h (\hat{u}_h^{(1)} - \hat{u}_h^{(2)}) & \text{B} & Z_h^* \end{cases} \quad (95)$$

$$\begin{cases} \hat{\sigma}_h^{(1)} - \hat{\sigma}_h^{(2)} = \hat{D}_h (\hat{\varepsilon}_h^{(1)} - \hat{\varepsilon}_h^{(2)}) & \text{B} & Z_h \end{cases} \quad (96)$$

$$\begin{cases} H_h (\hat{\sigma}_h^{(1)} - \hat{\sigma}_h^{(2)}) = 0 & \text{B} & V_h^* \end{cases} \quad (97)$$

Then, by virtue of (92)–(94), we have $(\hat{u}_h^{(1)}, \hat{\sigma}_h^{(1)}, \hat{\varepsilon}_h^{(1)}) \in (\hat{u}_h^{(2)}, \hat{\sigma}_h^{(2)}, \hat{\varepsilon}_h^{(2)})$. Therefore, the problem (48)–(50) is well-posed.

CONCLUSION

In conclusion, we note that condition (70), used in the consideration of the existence, uniqueness, and stability of the solution to the discrete problem (48)–(50), also plays an important role in proving the convergence of the mixed approximation process. At the same time, the lower bound estimates of the quantity 11111 are constructed for each specific projection-grid scheme and often are not a trivial task. Attempting to ignore condition (70) leads to an ill-conditioned system of equations upon mesh refinement, from which the discrete solution is determined with unsatisfactory accuracy.

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Article Title ("Work"):

Application of the Mixed Approximation for Solving Two

Problems of the Theory of Elasticity by the Finite Element Method

All Author(s):

Obid Abdullayev, Zokir Khudoyberdiyev, Khursand Ismoilov

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Title of Conference: **AMSMT2025**

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